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Toward a unified description of nuclei

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Nuclide Chart



Nuclide Chart



Nuclide Chart



Toward a unified description of nuclei

- <u>a global approach</u> which is applicable for the whole nuclear chart (light & heavy; neutron-deficient & neutron-rich)
- <u>a consistent treatment</u> for the ground-state and spectroscopic properties (low spin & high spin)
- <u>a unified description</u> of nuclear structure and nuclear reaction
- <u>a microscopic formalism</u> based directly on realistic nuclear interactions



State-of-the-art theories for nuclei



It would be interesting to investigate the intersections between different theories for a unified and comprehensive description of nuclei.

Density functional theory

The many-body problem is mapped onto an one-body problem



Figure from Drut PPNP 2010

$$E[\rho] \Rightarrow \hat{h} = \frac{\delta E}{\delta \rho} \Rightarrow \hat{h} \varphi_i = \varepsilon_i \varphi_i \Rightarrow \rho = \sum_{i=1}^A |\varphi_i|^2$$

The practical usefulness of the Kohn-Sham theory depends entirely on whether an Accurate Energy Density Functional can be found!

Density functional theory for nuclei

- ✓ The nuclear force is complicated
- ✓ More degrees of freedom: spin, isospin, pairing, ...
- ✓ Nuclei are self-bound systems
 DFT for the intrinsic density





- ✓ At present, all successful functionals are phenomenological not connected to any NN- or NNN-interaction
- ✓ Adjust to properties of nuclear matter and/or finite nuclei, and (in future) to ab-initio results



Relativistic vs Nonrelativistic

 \checkmark No relativistic kinematics necessary

$$\sqrt{p_F^2 + m_N^2} = m_N \sqrt{1 + 0.075}$$

✓ Large spin-orbit splitting



Similarity Renormalization Group

Relativistic Kohn-Sham Equation: Dirac Equation

$$\begin{pmatrix} m + \mathbf{V} + \mathbf{S} & \boldsymbol{\sigma} \cdot \boldsymbol{p} \\ \boldsymbol{\sigma} \cdot \boldsymbol{p} & -m + \mathbf{V} - \mathbf{S} \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix} = \varepsilon \begin{pmatrix} f \\ g \end{pmatrix}$$

$$\frac{dH(l)}{dl} = \begin{bmatrix} \eta(l), H(l) \end{bmatrix}, \quad \eta(l) = \begin{bmatrix} \beta, H(l) \end{bmatrix}$$

Bylev, Pirner, Phys. Lett. B 428, 329 (1998)

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Flow equation

$$\frac{dH(l)}{dl} = [\eta(l), H(l)], \quad \eta(l) = [\beta, H(l)]$$

Expansion with nucleon mass

$$\frac{1}{M}H(l) = \sum_{k=0}^{\infty} \frac{1}{M^k}H_k(l)$$

$$\blacksquare H(\infty) = \begin{pmatrix} H^F + M & 0 \\ 0 & H^D - M \end{pmatrix}$$

Bylev, Pirner, Phys. Lett. B 428, 329 (1998) Guo, Phys. Rev. C 85, 021302(R) (2012)

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Too slow convergence!

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Medium effects

$$\begin{pmatrix} m + \mathbf{V} + \mathbf{S} & \boldsymbol{\sigma} \cdot \boldsymbol{p} \\ \boldsymbol{\sigma} \cdot \boldsymbol{p} & -m + \mathbf{V} - \mathbf{S} \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix} = \varepsilon \begin{pmatrix} f \\ g \end{pmatrix}$$

Expansion with effective
nucleon mass
$$\tilde{M}(r) = M + S(r)$$

$$\frac{dH(l)}{dl} = \begin{bmatrix} \eta(l), H(l) \end{bmatrix}, \quad \eta(l) = \begin{bmatrix} \beta, H(l) \end{bmatrix} \implies H(\infty) = \begin{pmatrix} H^F + M & 0 \\ 0 & H^D - M \end{pmatrix}$$

Ren, PWZ, Phys. Rev. C, 100, 044322 (2019)

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Toward a bridge between relativistic and nonrelativistic DFT

Ren, PWZ, Phys. Rev. C, 102, 021301(R) (2020) Editors' Suggestion

Covariant density functional: PC-PK1

~10 parameters fitted to 60 spherical nuclei ...

rms-deviation 2.96 MeV 1.14 PC-PK1 DD-PC1 ТМА DD-MEδ DD-ME2 NL3*

> Agbemava PRC 2014 Geng PTP 2005

Best density-functional description for nuclear masses so far!

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Outline

Nuclear landscape

(light and heavy nuclei; neutron-deficient and neutron-rich nuclei)

- Nuclear spectroscopy (ground and excited states; low- and high- angular momentum states)
- Nuclear dynamics

(nuclear structure and reaction)

Nuclear interactions

(towards a microscopic formalism based on bare nuclear interactions)

• Summary

Nuclear landscape

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Neutron Number N

What will the new facilities bring us?

More isotopes ...

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Magnetic and Anti-magnetic Rotation

Tilted axis cranking CDFT

PWZ, Peng, Liang, Ring, Meng, PRL 107, 122501 (2011) PWZ, Zhang, Peng, Liang, Ring, Meng, PLB 699, 181 (2011)

Experimental observations in 2007: Energy Spectrum

Joshi, et al., PRL 98, 102501 (2007)

A pair of strongly coupled bands observed Chiral bands? But why crossing?

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A pair of strongly coupled bands observed Chiral bands? But why crossing?

Experimental observations in 2014: Transition strength

Rather, et al., PRL 112, 202503 (2014)

Lieder, et al., PRL 112, 202502 (2014)

B(M1) and B(E2) values were measured in both experimentsA third band is reported in Lieder's experimentChiral bands ? But why crossing ? Why three bands ?

TAC-CDFT calculations

2qp

PWZ PLB 773,1 (2017) PWZ, Wang, Chen PRC 99, 054319 (2019)

(C)DFT and Shell Model

(C)DFT

Shell Model

Universal density functionals

Symmetry broken Single config. fruitful physics No Configuration mixing

Applicable for almost all nuclei
 No spectroscopic properties

Non-universal effective interactions

No symmetry broken Single config. little physics Configuration mixing

intractable for deformed heavy nuclei
 spectroscopy from multi config.

a theory combining the advantages from both approaches

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Configuration Interaction Projected DFT (CI-PDFT)

PWZ, Ring, Meng, PRC 94 (2016) 041301(R)

 $0\nu\beta\beta$?

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Next step: time-odd interaction; beyond 2-qp configurations;

Towards neutron-rich nuclei?

Towards triaxial nuclei ?

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Time-dependent density functional theory

The many-body problem is mapped onto a one-body problem

Runge-Gross Theorem

There is a unique mapping between the time dependent external potential and the density, for many body systems evolving from a given initial state.

Time-dependent Kohn-Sham DFT

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Runge and Gross, PRL 52, 997 (1984)

$$i\hbar \frac{\partial \phi_i(\boldsymbol{r},t)}{\partial t} = \left[-\frac{\hbar^2}{2m} \nabla^2 + v_{KS}[\rho(\boldsymbol{r},t)] \right] \phi_i(\boldsymbol{r},t) \qquad \rho(\boldsymbol{r},t) = \sum_i^N |\phi_i(\boldsymbol{r},t)|^2$$

Resonant scattering of ⁴He + ⁸Be

The linear-chain states are generated, and then evolve to a triangular configuration, and finally to a more compact shape.

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Adding more neutrons: ⁴He +¹⁰Be

Ren, PWZ, Meng, PLB 801, 135194 (2020)

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✓ The metastable linear chains can be formed in ⁴He+⁸Be and ⁴He+¹⁰Be collisions.

✓ During the time evolution of the linear-chain configuration, moving clusters can be found.

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Ren, PWZ, Meng, PLB 801, 135194 (2020)

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Statical calculations

Rod-like shapes could be realized in nuclei with large total spin and isospin.

PWZ, Itagaki, Meng, PRL 115, 022501 (2015)

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Rod-like shapes could be realized in nuclei with large total spin and isospin.

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Exp @RIKEN

where \Im is the moment of inertia of the nucleus. The linearity allows us to interpret the levels as a rotational band, and the low $\hbar^2/2\Im = 0.19$ MeV implies the nucleus could be strongly deformed, consistent with the interpretation of an LCCS. Although we ob-

Yamaguchi et al., PLB 766 (2017) 11-16

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Complicated nuclear force

Short-range repulsive; Three-body force; Covariance

Neutron Drops

.

Why study neutron drops? (A pure simple toy model?)

Pudliner, et al., PRL 76, 2416 (1996) Gandolfi, et al., PRL 106, 012501 (2011) Maris, et al., PRC 87, 054318 (2013) Potter, et al., PLB 739, 445 (2014) PWZ and Gandolfi, PRC 94 (2016) 041302(R) Shen, et al., PLB 778, 344 (2018)

- 1. benchmark various nuclear many-body methods
- 2. calibrate nuclear interactions / energy density functionals
- 3. model neutron-rich nuclei and neutron stars
- 4. predict few neutron resonances
- 5. connections with the skin thickness and symmetry energy

Neutron drop as a "toy" neutron-rich nucleus

Neutron skin thickness

PWZ and Gandolfi PRC 94 (2016) 041302(R)

Constraining three-body forces

Radii: QMC calculations with various Hamiltonians, compared to what is extracted from "experiments". PWZ and Gandolfi PRC 94 (2016) 041302(R)

With the development of the high-accuracy measurements of neutron skin thickness, the radii obtained for neutron drops will <u>provide a useful constraint</u> for realistic three neutron forces.

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Neutron drop as a "toy" neutron star

Going to high densities ...

Tong, PWZ, Meng, PRC 101, 035802 (2020)

The strong correlation can be used to constrain the radius of the neutron drop

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Summary

Covariant density functional theory is being <u>improved and extended</u> for a unified description of nuclei.

- Nuclear landscape
 a systematic study with PC-PK1
 triaxiality + dynamical correlations!
- Nuclear spectroscopy a new tool for global studies of nuclear spectroscopy: CI-PDFT merits of (C)DFT and Shell Model preserved
- Nuclear dynamics time-dependent CDFT has been developed pave the way for a unified description of nuclear structure and reaction
- Nuclear interaction

neutron drop, neutron star, and neutron skin are found to be strongly correlated likely to have an enduring impact on the understanding of multi-neutron interactions

Collaborations

Beijing Jie Meng Jing Peng Zhengxue Ren Yakun Wang Shuangquan Zhang Munich Qibo Chen Peter Ring

Chongqing Zhipan Li **Kyoto** Naoyuki Itagaki

Los Alamos Stefano Gandolfi

Thank you for your attention!

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